# Relativistic Projectile Coulomb Excitation to the Yrast and Non-yrast 2+ States of ${ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$ 

Take R. Saito

Gesellschaft für Schwerionenforschung (GSI),
A research center of Helmholtz Association and
Johannes Gutenberg-Universität Mainz
for the RISING collaboration


## Physics motivation

- Nuclear phase transition around A~130



## Physics motivation

Triaxiality and $\gamma$-softness play a role in chiral structures

Indication of triaxiality and $\gamma$-softness near the ground state

- Lower $2^{+}{ }_{2}$ energy and large $\mathcal{B}\left(E 2 ; 2^{+}{ }_{2} \rightarrow 0^{+}\right)$
- $\boldsymbol{\gamma}$-softness
- Lower $2^{+}{ }_{2}$ energy than $4^{+}{ }_{1}$ and large $\mathcal{B}\left(E 2 ; 2^{+}{ }_{2} \rightarrow 2^{+}{ }_{1}\right)$
- Triaxiality
- Large scale microscopic Monte Carlo Shell Model Calculation is available
- Reproduced nuclear properties of neutron rich $\mathcal{B a}$ is otopes


## Goal of the experiment

- To measure $B(E 2)$ of transitions depopulating $\mathbf{2 ~}^{+}{ }_{1}$ and $2^{+}$, states in ${ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$
- Relativistic Coulomb excitation of secondary ${ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$ projectiles with FRS at GSI
- $\gamma$-ray measurements with RISING Ge detector array

|  | $\begin{gathered} \mathbf{N d} 135_{12.4 \mathrm{~m}}^{1.2(-)} \\ \mathrm{EC} \\ \hline \end{gathered}$ |  | $\begin{aligned} & \mathrm{Nd137} \\ & 38.5 \mathrm{~m} \\ & 1 / 2+ \\ & c \quad * \\ & \hline \end{aligned}$ |  | $\begin{gathered} \mathrm{Nd} 139_{29.7 \mathrm{~m}}^{3 / 2+} \\ \mathrm{EC}^{3} \end{gathered}$ | $\begin{aligned} & \hline \begin{array}{c} \text { Nd140 } \\ 3.37 \mathrm{~d} \\ \mathbf{0}+ \\ \mathrm{EC} \end{array} \\ & \hline \end{aligned}$ | $\begin{aligned} & \mathrm{Nd} 141_{2.49 \mathrm{~h}} \\ & 3 / 2+ \\ & \mathrm{EC} \\ & \hline \end{aligned}$ | $\begin{gathered} \mathrm{Nd} 142 \\ 0+ \\ 27.13 \end{gathered}$ | $\begin{gathered} \hline \text { Nd143 } \\ 7 / 2- \\ 12.18 \end{gathered}$ | $\begin{gathered} \mathrm{Nd144} \\ 2.29 \mathrm{E}+15 \mathrm{y} \\ 0+ \\ \& \quad 23.80 \end{gathered}$ | $\begin{gathered} \mathrm{Nd} 145 \\ 7 / 2- \\ 8.30 \end{gathered}$ | $\begin{gathered} \mathrm{Nd} 146 \\ 0+ \\ 17.19 \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\begin{aligned} & \hline \operatorname{Pr} 133 \\ & 6.5 \mathrm{~m} \\ & (3 / 2+) \\ & \text { EC } \end{aligned}$ | Pr1 17 m $2-$ $2-$ EC |  | $\underbrace{}_{\substack{\text { Pr136 } \\ \text { E. } \\ \text { 2+ m }}}$ | $\begin{gathered} \text { Pr13 } \\ 1.281 \\ 5 / 2+ \\ \mathbf{E C} \end{gathered}$ | $\underbrace{\substack{1.45 \mathrm{~m} \\ 1+}}_{\mathrm{EC}^{\text {Pr138 }}} \quad$ * | Pr139 4.41 h $5 / 2+$ EC | Pr140 <br> $\substack{3.39 \mathrm{~m} \\ 1+\\ \text { EC }}$ | Pr14 | $\begin{gathered} \text { Pr142 } \\ 19.12 \mathrm{~h} \\ 2-\quad * \end{gathered}$ |  | $\begin{gathered} \operatorname{Pr} 144 \\ 17.28 \mathrm{~m} \\ 0- \end{gathered}$ | Pr145 5.984 h $7 / 2+$ $3-$ |
| $\begin{aligned} & \hline \text { Ce132 } \\ & \begin{array}{c} 3.51 \mathrm{~h} \\ \mathbf{0}^{+} \\ \text {EC }^{+} \end{array} \quad * \end{aligned}$ | Ce133 $\mathrm{EC}^{97 \mathrm{~m}}$ $1 / 2+$ |  | $\begin{gathered} \text { Ce135 } \\ 17.7 \mathrm{~h} \\ 1 /(+) \end{gathered} \mathrm{c}^{2}$ | $\begin{gathered} \text { Ce13 } \\ 0+ \\ 0.19 \end{gathered}$ | $\begin{gathered} \text { Ce137 } \\ 9.0 \mathrm{~h} \\ 3 / 2+ \\ \text { EC } \end{gathered} *$ | Ce13 $0+$ 0.25 | $\begin{aligned} & \text { Ce13 } \\ & \text { 137.641 } \\ & 3 / 2+ \\ & \text { EC } \\ & \hline \end{aligned}$ | Ce14 <br> 0+ <br> 88,48 | $\begin{gathered} \text { Ce141 } \\ \text { 32.501 } \\ 7 / 2- \end{gathered}$ | $\begin{gathered} \text { Ce142 } \\ \begin{array}{c} 5 \mathrm{E}+16 \mathrm{y} \\ 0+ \\ 11.08 \end{array} \\ \hline \end{gathered}$ | $\begin{gathered} 33.039 \mathrm{~h} \\ 3 / 2- \end{gathered}$ | $\begin{aligned} & \hline \mathbf{C e 1 4 4} \\ & 284.893 \mathrm{~d} \\ & 0+ \\ & \beta- \\ & \hline \end{aligned}$ |
| $\begin{array}{r} \text { La13 } \\ 59 \mathrm{~m} \\ 3 / 2+ \\ \mathrm{EC} \\ \hline \end{array}$ |  | $\begin{gathered} \text { La13 } \\ 3.912 \\ 5 / 2+ \\ \text { EC } \end{gathered}$ | $\begin{gathered} \text { La13 } \\ \mathbf{6 . 4 5} \mathrm{n} \\ 1+ \\ \text { EC } \end{gathered}$ | $\begin{gathered} \text { La13s } \\ 19.5 \mathrm{~h} \\ 5 / 2+ \\ \mathbf{E C} \end{gathered}$ |  | $\begin{array}{\|c} \hline \mathbf{L a 1 3 7} \\ 6 \mathrm{EA} \mathbf{y} \\ 7 / 2+ \\ \text { EC } \\ \hline \end{array}$ |  | La13 7/2+ 99,909 | $\begin{gathered} \text { La140 } \\ 1.6781 \mathrm{~d} \\ 3- \end{gathered}$ | $\substack{\text { La141 } \\ 3.92 \mathrm{~h} \\ (72 \mathrm{~h}) \\ \beta \\ \beta}$ | $\begin{gathered} \text { La142 } \\ \substack{91.1 \mathrm{~m} \\ 2-\\ \beta} \\ \hline \end{gathered}$ | $\begin{aligned} & 14.2 \mathrm{~m} \\ & (7 / 2)+ \\ & \beta- \end{aligned}$ |
| $\begin{gathered} \text { Ba130 } \\ 0+ \\ 0.106 \end{gathered}$ | $\begin{aligned} & \text { Ba131 } \\ & 11.50 \mathrm{~d} \\ & 1 / 2+ \\ & \mathrm{EC}^{*} * \end{aligned}$ | Ba13 $0+$ 0.101 | $\begin{gathered} \begin{array}{c} \text { Ba133 } \\ 10.51 \mathrm{y} \\ 1 / 2+ \\ \text { EC } \end{array} \quad * \cdot \end{gathered}$ | Ba1 | 3913 $3 / 2+$ 6.592 | Ba13 $0+$ 7.854 | Ba13 $3 / 2+$ 11.23 | Ba13 $0+$ 71.70 | $\begin{gathered} 83.06 \mathrm{~m} \\ 7 / 2- \end{gathered}$ |  |  | $\underset{\substack{10.6 \mathrm{~m} \\ 0+}}{\substack{\mathrm{Ba} \\ \hline}}$ |
| $\begin{aligned} & \text { Cs129 } \\ & 32.06 \mathrm{~h} \\ & 1 / 2+ \\ & \mathrm{EC} \\ & \hline \end{aligned}$ | $\begin{gathered} \text { Cs130 } \\ 29.21 \mathrm{~m} \\ 1+ \\ \text { EC }, \beta \end{gathered} \quad *$ | $\begin{gathered} \text { Cs131 } \\ 9.689 \mathrm{~d} \\ 5 / 2+ \\ \mathrm{EC} \\ \hline \end{gathered}$ | $\begin{gathered} \text { Cs132 } \\ 6.479 \mathrm{~d} \\ 2+ \\ \mathbf{E C}, \beta \end{gathered}$ | Cs13 7/2+ 100 | $\begin{gathered} \text { CS1 } \\ 2.0648 \\ 4+ \\ \mathbf{E C}, \beta \end{gathered}$ | $\begin{gathered} \text { Cs135 } \\ 2.3 \mathrm{E}+6 \\ 7 / 2+ \\ \beta \\ \hline \end{gathered}$ | $\int_{3}^{13.16 \mathrm{~d}} 5$ | $\begin{gathered} 30.07 \mathrm{y} \\ 7 / 2+ \end{gathered}$ | $\underbrace{33.41 \mathrm{~m}}_{3} \quad 3$ |  |  |  |
| $\begin{gathered} \text { Xe128 } \\ 0+ \\ 1.91 \end{gathered}$ | $\begin{gathered} \text { Xe129 } \\ 1 / 2+ \\ 26.4 \end{gathered}$ | Xe13 | Xe131 $3 / 2+$ 21.2 | Xe132 | $\begin{gathered} \text { Xe133 } \\ 5.243 \mathrm{~d} \\ 3 / 2+ \\ 2 \end{gathered}$ | $\begin{gathered} \text { Xe134 } \\ 0+ \\ 10.4 \\ \hline \end{gathered}$ | $\begin{gathered} \text { Xe135 } \\ \substack{9.14 \mathrm{~h} \\ 3 / 2+\\ *} \end{gathered}$ | $\begin{gathered} 2.3021 \mathrm{y} \\ 0+ \end{gathered}$ | $\begin{gathered} \mathbf{N e l s i n}_{7 / 2-} \end{gathered}$ | $\begin{gathered} 14.08 \mathrm{~m} \\ 0+ \end{gathered}$ | $\begin{gathered} \text { Ae13y } \\ 39.68 \mathrm{~s} \\ 3 / 2- \end{gathered}$ | $\begin{array}{\|c} \hline \begin{array}{c} \text { Xe140 } \\ 13.60 \mathrm{~s} \\ 0+ \\ \beta- \\ \beta \end{array} \\ \hline \end{array}$ |

## Fragment separator



## RISING setup at the final focal plane



## Doppler shift correction



## Selection of Coulomb excitation events: 





[^0]
## Gamma-ray spectra





## Particle $-\gamma$ angular correlation

- Correlation on $\boldsymbol{\theta}$
- Rising-wise sorting at the rest frame
- Isotropic distribution observed
- Can not reproduced by calculations with alignments
- Efficiency calibration being crucial

- Correlation on $\varphi$ difference between $\gamma$-rays and outgoing particles, event-by-event
- Using all Ge-CATE phase space
- Efficiency calibration cancelled
- Isotropic distribution observed
- Can not reproduced by calculations with alignments


Surprisingly, no (or very small) alignment was observed

## Deducing $B(E 2)$

|  | ${ }^{134} \mathrm{Ce}$ on ${ }^{197} \mathrm{Au}$ at 126 A MeV |  |  | ${ }^{136} \mathrm{Nd}$ on ${ }^{197} \mathrm{Au}$ at 126 A MeV |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $\mathrm{E}_{\gamma}[\mathrm{keV}]$ | $\mathrm{N}_{\gamma}$ | $\varepsilon[\%]$ | $\mathrm{E}_{\gamma}[\mathrm{keV}]$ | $\mathrm{N}_{\gamma}$ | $\varepsilon[\%]$ |
| $2^{+}{ }_{1} \rightarrow 0^{+}$ | $4091713 \pm 101$ |  | $1.98 \pm 0.06$ | 374 | $3039 \pm 130$ | $1.98 \pm 0.06$ |
| $2^{+}{ }_{2} \rightarrow 0^{+}$ | 966 |  | $1.51 \pm 0.05$ | 862 | $156 \pm 49$ | $1.58 \pm 0.05$ |
| $2^{+}{ }_{2} \rightarrow 2^{+}{ }_{1}$ | 557 | <149 | $1.85 \pm 0.05$ | 489 | $183 \pm 56$ | $1.91 \pm 0.06$ |
|  | $\begin{gathered} \mathrm{N}_{\text {pro }} \\ 1895843 \pm 1377 \end{gathered}$ | DAQ livetime [\%] 77 |  | $\begin{gathered} \mathrm{N}_{\text {pro }} \\ 1986411 \pm 1409 \end{gathered}$ | DAQ livetime [\%] 79 |  |

Known value: ${ }^{134} \mathrm{Ce} 2^{+}{ }_{1} \rightarrow 0^{+} \mathrm{B}(\mathrm{E} 2), 52$ (5) W.u.; reference
Absolute efficiency with Lorentz boost with $\beta=0.42$ and DAQ livetime taken into account

- Isotropic distribution at the rest frame assumed
- Gamma-ray intensities normalized by particle numbers on the target


## $B(E 2)$ values in W.u.

|  | ${ }^{132} \mathrm{Ba}$ | ${ }^{134} \mathrm{Ce}$ | ${ }^{136} \mathrm{Nd}$ |
| :---: | :---: | :---: | :---: |
| $2^{+}{ }_{1} \rightarrow 0^{+}$ | 43 (4) | 52 (5) | 80 (11) |
|  |  | 77 (26) | 97 (27) |
| $2^{+}{ }_{2} \rightarrow 0^{+}$ | 3.9 (4) | < 11 | 11 (3) |
|  |  |  | 13 (5) |
| $2^{+}{ }_{2} \rightarrow 2^{+}{ }_{1}$ | 144 (14) | < 140 | 182 (93) |

- Complete measurements for ${ }^{136} \mathrm{Nd}$
- Only upper limit given to $2^{+}{ }_{2} \rightarrow \mathrm{O}^{+}$and $2^{+}{ }_{2} \rightarrow \mathrm{O}^{+}$in ${ }^{134} \mathrm{Ce}$

B(E2) values with normalization to the target ${ }^{197}$ Au Coulomb excitation used for cross-checking

## Comparison to the theoretical calculations

- Large scale microscopic Monte Carlo Shell Model Calculation

Still in progress by Otsuka of Tokyo University

- Naive Macroscopic Asymmetric Rotor model calculations
- Experimental information on the transition probability is limited only up to the $2^{+}$atates


## Asymmetric Rotor Model (ARM) for ${ }^{132} \mathrm{Ba},{ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$


$\frac{B\left(E 2 ; 2_{2} \rightarrow 2_{1}\right)}{B\left(E 2 ; 2_{2} \rightarrow 0\right)}=\frac{20}{7} \frac{\sin ^{2}(3 \gamma)}{\sqrt{9-8 \sin ^{2}(3 \gamma)}-3+2 \sin ^{2}(3 \gamma)}$

- ${ }^{134}$ Ce: from the known branching ratio
- ${ }^{136} \mathcal{N} d$ : from this experiment


## Asymmetric Rotor Model (ARM) for ${ }^{132} \mathrm{Ba},{ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$

- $\gamma$-rigid ARM
A.S. Davydov and G.F. Filippov, Nucle. Pfys. 8, 237(1958)
- $\gamma$-soft ARM with $\mu=0.5$
$\mathcal{A} . S . \operatorname{Davydov}$ and $\mathcal{A} . \mathcal{A}$. Chaban, Nucle. Pfys. 20, 499 (1960)



## Summary

- Relativistic projectile Coulomb excitation with RI-beams of ${ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$ to $2^{+}{ }_{1}$ and $2^{+}{ }_{2}$ with FRS-RISING at GSI
- The first relativistic Coulomb excitation to non-yrast states
- No alignment observed
- Relative $B(E 2)$ measurement normalized to the known $B(E 2)$ of $\mathbf{2}^{+}{ }_{1} \rightarrow 0^{+}$in ${ }^{134} \mathrm{Ce}$
- Data compared to asymmetric rotor model
- Indication of a triaxial soft rotor
- Microscopic calculations with Monte Carlo Shell Model are in progress


## Backup slides

## Deducing $B(E 2)$

|  | ${ }^{134} \mathrm{Ce}+{ }^{197} \mathrm{Au}$ at 126 A MeV |  |  | ${ }^{136} \mathrm{Nd}+{ }^{197} \mathrm{Au}$ at 126 A MeV |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Transitions | $\mathrm{E}_{\gamma}[\mathrm{keV}]$ | $\mathrm{N}_{\gamma}$ | $\varepsilon[\%](\beta=0.42)$ | $\mathrm{E}_{\boldsymbol{\gamma}}[\mathrm{keV}]$ | $\mathrm{N}_{\gamma}$ | $\varepsilon[\%](\beta=0.42)$ |
| $2_{1}^{+} \rightarrow 0^{+}$ | 409.2 | $1713 \pm 101$ | $1.98 \pm 0.06$ | 373.7 | $3039 \pm 130$ | $1.99 \pm 0.06$ |
| $2_{2}^{+} \rightarrow 0^{+}$ | 965.7 |  | $1.51 \pm 0.05$ | 862.4 | $156 \pm 49$ | $1.58 \pm 0.05$ |
| $2_{2}^{+} \rightarrow 2_{1}^{+}$ | 556.6 | < 149 | $1.85 \pm 0.06$ | 488.6 | $183 \pm 56$ | $1.91 \pm 0.06$ |
| ${ }^{197} \mathrm{Au}$ transitions$7 / 2^{+} \rightarrow 3 / 2^{+}$ | $\begin{gathered} \mathrm{E}_{\gamma}[\mathrm{keV}] \\ 547.5 \\ \hline \end{gathered}$ | $\begin{gathered} \mathrm{N}_{\gamma} \\ 130 \pm 41 \\ \hline \end{gathered}$ | $\begin{gathered} \varepsilon[\%](\beta=0.0) \\ 0.97 \pm 0.03 \end{gathered}$ | $\begin{gathered} \mathrm{E}_{\gamma}[\mathrm{keV}] \\ 547.5 \\ \hline \end{gathered}$ | $\begin{gathered} \mathrm{N}_{\gamma} \\ 171 \pm 44 \\ \hline \end{gathered}$ | $\begin{gathered} \varepsilon[\%](\beta=0.0) \\ 0.97 \pm 0.03 \\ \hline \end{gathered}$ |
|  | $\begin{gathered} \mathrm{N}_{\text {pro }} \\ 189583 \pm 1377 \end{gathered}$ |  | $\begin{aligned} & \text { DAQ livetime [\%] } \\ & 77 \end{aligned}$ | $\begin{gathered} \mathrm{N}_{\text {pro }} \\ 1986411 \pm 1409 \end{gathered}$ |  | $\begin{gathered} \text { DAQ livetime [\%] } \\ 79 \end{gathered}$ |

Known value: ${ }^{134} \mathrm{Ce} 2^{+}{ }_{1} \rightarrow 0^{+} \mathrm{B}(\mathrm{E} 2), 52$ (5) W.u.; reference

- Absolute efficiency with Lorentz boost with $\beta=0.42$ and DAQ livetime taken into account
- Isotropic distribution at the rest frame assumed
- Gamma-ray intensities normalized by particle numbers on the target


## Absolute efficiency calibration

- $\gamma$ - $\gamma$ coincidence measurement with ${ }^{60} \mathrm{Co} \gamma$-ray source
- Absolute efficiency for 1.173 and $1.333 \mathcal{M e} \mathcal{V} \gamma$-rays
- Calibration with an assumption of isotropic $\gamma-\gamma$ angular correlation
- Actual $\gamma$ - $\gamma$ angular correlation gives difference of $1 \%$ in the calibration
$\rightarrow$ systematic error
- Relative efficiency calibration (energy dependence) with ${ }^{152} \mathrm{Eu}$ $\gamma$-ray source
- Lorentz-boost to the rest frame



## Selection of Coulomb excitation events: $\gamma$-multiplicity




## Improved particle tracking



- Old tracking method: MW41 - MW42 - CATE
- Scattering in $\mathcal{M A}$ IS $->$ wrong vertex reconstruction
- New method with target-Si detector
- $\mathcal{N}$ o affection from $\mathcal{M A}$ IS IC on the tracking
- Easier vertex reconstruction


## Asymmetric Rotor Model (ARM) for ${ }^{132} \mathrm{Ba},{ }^{134} \mathrm{Ce}$ and ${ }^{136} \mathrm{Nd}$

- Very naive calculations
- For transitions from $\mathbf{2 ~}_{1}$
- ${ }^{132} \mathcal{B a}: \beta=0.19$
- ${ }^{134} \mathrm{Ce}: \beta=0.20$
- ${ }^{136}$ ㄱ $d: \beta=0.24$
- For transitions from $\mathbf{2 +}_{2}$

$$
\begin{aligned}
& B\left(E 2 ; 2_{1}^{+} \rightarrow 0^{+}\right)=\frac{e^{2} Q_{0}^{2}}{16 \cdot \pi} \cdot \frac{1}{2} \cdot\left[1+\frac{3-2 \cdot \sin ^{2}(3 \gamma)}{\sqrt{9-8 \cdot \sin ^{2}(3 \gamma)}}\right] \\
& Q_{0}=\frac{3 \cdot Z \cdot R^{2} \cdot \beta}{\sqrt{5 \pi}}
\end{aligned}
$$

$$
\begin{gathered}
B\left(E 2 ; 2_{2}^{+} \rightarrow 0^{+}\right)=\frac{e^{2} Q_{0}^{2}}{16 \cdot \pi} \cdot \frac{1}{2} \cdot\left[1-\frac{3-2 \cdot \sin ^{2}(3 \gamma)}{\sqrt{9-8 \cdot \sin ^{2}(3 \gamma)}}\right] \\
B\left(E 2 ; 2_{2}^{+} \rightarrow 2_{1}^{+}\right)=\frac{e^{2} Q_{0}^{2}}{16 \cdot \pi} \cdot \frac{10}{7} \cdot \frac{\sin ^{2}(3 \gamma)}{\sqrt{9-8 \cdot \sin ^{2}(3 \gamma)}} \\
\frac{B\left(E 2 ; 2_{2} \rightarrow 2_{1}\right)}{B\left(E 2 ; 2_{2} \rightarrow 0\right)}=\frac{20}{7} \frac{\sin ^{2}(3 \gamma)}{\sqrt{9-8 \sin ^{2}(3 \gamma)}-3+2 \sin ^{2}(3 \gamma \gamma}
\end{gathered}
$$

## EUROBALL cluster detectors



- 15 Euroball cluster Ge detectors (105 crystals)
- Energy resolution for 1 MeV $\gamma$-ray with 126 A MeV RIbeams: ~2.0 \%

|  | Cluster <br> detector | Angle | Target <br> distance |
| :--- | :---: | :---: | :---: |
| Ring \#1 | 5 | $15.9^{\circ}$ | $\sim 720 \mathrm{~mm}$ |
| Ring \#2 | 5 | $33.0^{\circ}$ | $\sim 720 \mathrm{~mm}$ |
| Ring \#3 | 5 | $36.0^{\circ}$ | $\sim 720 \mathrm{~mm}$ <br> mm |


[^0]:    Cut with $0.8 \sim 1.8$ degrees is optimum

